

Fig. 5. Temperature dependence of magnetization of Mn_3Ge_2 alloy at atmospheric pressure (solid curves) and at 9700 atm (dashed curves).

TABLE 1. Thermodynamic Data Characterizing the Low-Temperature Transition in ${\rm Mn}_3{\rm Ge}_2$

θ ₁ , °K	Δy, cm ³ /g	Δσ ₈ , G•cm ³ /g	AS, ergs/g·deg	ΔQ, cal/g
158	-1.55 · 10-5	2.0	0.8 · 10 ⁴ 5.0 · 10 ⁴	0.030 0.190

 $\sigma = \sigma_S + \chi H$; extrapolation of the obtained lines to a zero field made it possible to determined σ_S . These values were then corrected for the Mn_3Ge_2 phase content since, as noted above, our samples were an eutectic of Mn_3Ge_2 and Ge. The change in specific volume ΔV has been found from dilatometric data and the sample density $\rho = 6.44 \ g/cm^3$ measured by hydrostatic weighing.

As seen in Table 1 the change in entropy $\Delta S_1 = 0.8 \cdot 10^4 \, \mathrm{ergs/g} \cdot \mathrm{deg} \, \mathrm{calculated} \, \mathrm{from} \, \mathrm{magnetic} \, \mathrm{measurements} \, \mathrm{differs} \, \mathrm{considerably} \, \mathrm{from} \, \Delta S_2 = 5.0 \cdot 10^4 \, \mathrm{ergs/g} \cdot \mathrm{deg} \, \mathrm{found} \, \mathrm{from} \, \mathrm{the} \, \mathrm{shift} \, \mathrm{of} \, \Theta_1 \, \mathrm{with} \, \mathrm{pressure} \, \mathrm{and} \, \mathrm{from} \, \mathrm{the} \, \mathrm{change} \, \mathrm{in} \, \mathrm{volume} \, \mathrm{at} \, \mathrm{the} \, \mathrm{point} \, \mathrm{of} \, \mathrm{transition}. \, \mathrm{Consequently,} \, \mathrm{the} \, \mathrm{obtained} \, \mathrm{data} \, \mathrm{are} \, \mathrm{suitable} \, \mathrm{only} \, \mathrm{for} \, \mathrm{a} \, \mathrm{qualitative} \, \mathrm{comparison} \, \mathrm{with} \, \mathrm{the} \, \mathrm{Kittel} \, \mathrm{theory, which} \, \mathrm{is} \, \mathrm{based} \, \mathrm{on} \, \mathrm{the} \, \mathrm{exchange-in-version} \, \mathrm{mechanism} \, [8].$

This theory states that the change of magnetic transition temperature with pressure depends on Young's modulus and on the thermal expansion coefficient in the paramagnetic temperature intervals:

$$\frac{d\theta}{dP} = \frac{1}{E\alpha_{\rm p}}$$
.

Our experimental data give Young's modulus of $\mathrm{Mn_3Ge_2}$ as $5.50 \cdot 10^{11} \, \mathrm{dyn/cm^2}$ whereas the Kittel equation gives $\mathrm{E} = 2.5 \cdot 10^{14} \, \mathrm{dyn/cm^2}$; the sign of $\mathrm{d}\Theta_1/\mathrm{dP}$ also does not agree with theoretical conclusions. As was already mentioned, the magnetic transition in $\mathrm{Mn_3Ge_2}$ which takes place at the point Θ_1 with rising remperature is accompanied by constriction of the crystal lattice, whereas the theory [10] predicts lattice expansion in the case of an AF \rightarrow F transition. The change in lattice parameter in the AF \rightarrow F transition is determined by

$$\Delta a = a_{\mathrm{F}} - a_{\mathrm{AF}} = \frac{2\rho}{R} M^2,$$

where ρ is the rate of change of the exchange interaction as a function of interatomic spacing, $R=E/\alpha^2$, and M is the sublattice magnetization. This expression makes it clear that the sign of Δa is governed by the sign of ρ , i.e., by the sign of the derivative of the change of magnetic transition temperature with pressure. It should be mentioned in this connection that the negative sign of the $d\Theta_1/dH$ effect observed experimentally also does not agree with the Kittel expression

$$\frac{d\theta}{dH} = -\frac{1}{\rho M} \left(\frac{\partial a}{\partial T}\right)_{P}.$$

The exchange-inversion theory of C. Kittel has been further expanded in [14]. The entropy change